

# Shocks power Tidal Disruption Events

Ryu, Krolik, Piran, Noble, Avara

arXiv: 2505.05333

Reported by Shiyan Zhong

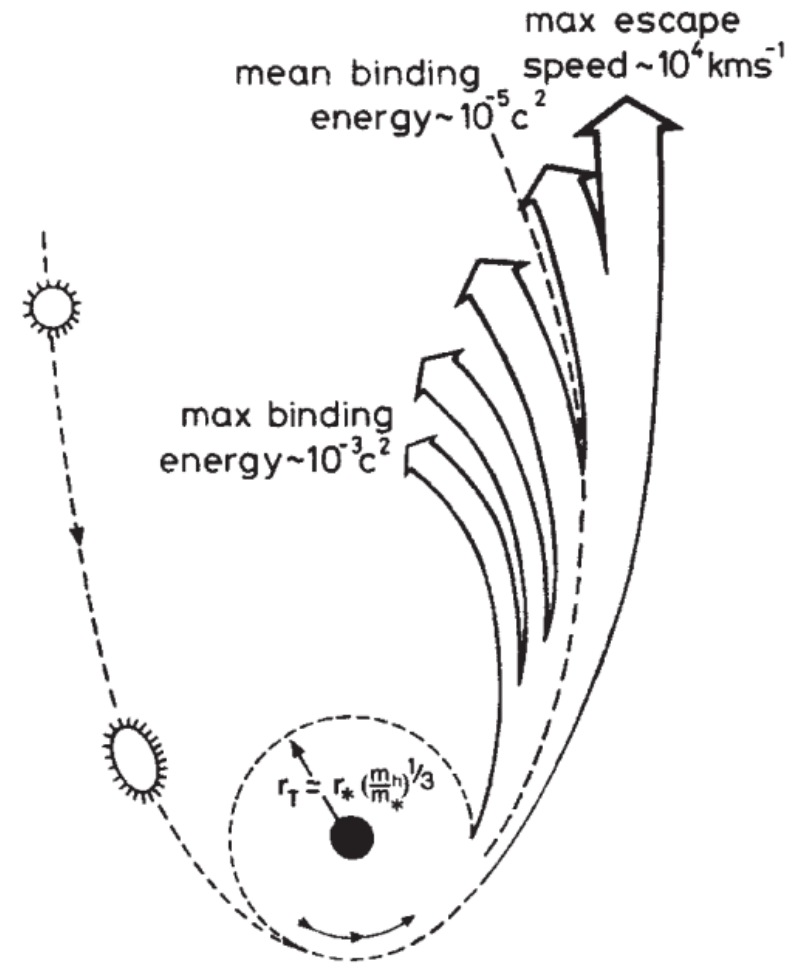
# Introduction

- What is TDE?

A star approaching to an MBH is torn apart and accreted, producing a luminous flare that can last for months to years.

Classic picture (Rees 1988):

- half of the debris are bound to the SMBH
- the debris quickly circularize to an accretion disk
- Small accretion disk (size  $\sim 2r_p$ )
- High temperature  $\rightarrow$  soft X-ray, EUV
- Light curve closely follow the mass fall back rate  $\rightarrow$  Super Eddington accretion



Rees (1988)

# Introduction

- What is TDE?

A star approaching to an MBH is torn apart and accreted, producing a luminous flare that can last for months to years.

Classic picture (Rees 1988):

- half of the debris are bound to the SMBH
- the debris quickly circularize to an accretion disk
- Small accretion disk (size  $\sim 2r_p$ )
- High temperature  $\rightarrow$  soft X-ray, EUV
- Light curve closely follow the mass fall back rate  $\rightarrow$  Super Eddington accretion

Challenged by observation:

- Observed luminosity rarely exceed Eddington
- Many TDEs are discovered in UV/Optical surveys
- Effective temperature is low: few  $10^4$  K

**Inverse Energy Crisis** (Piran+2015): more energy than demand

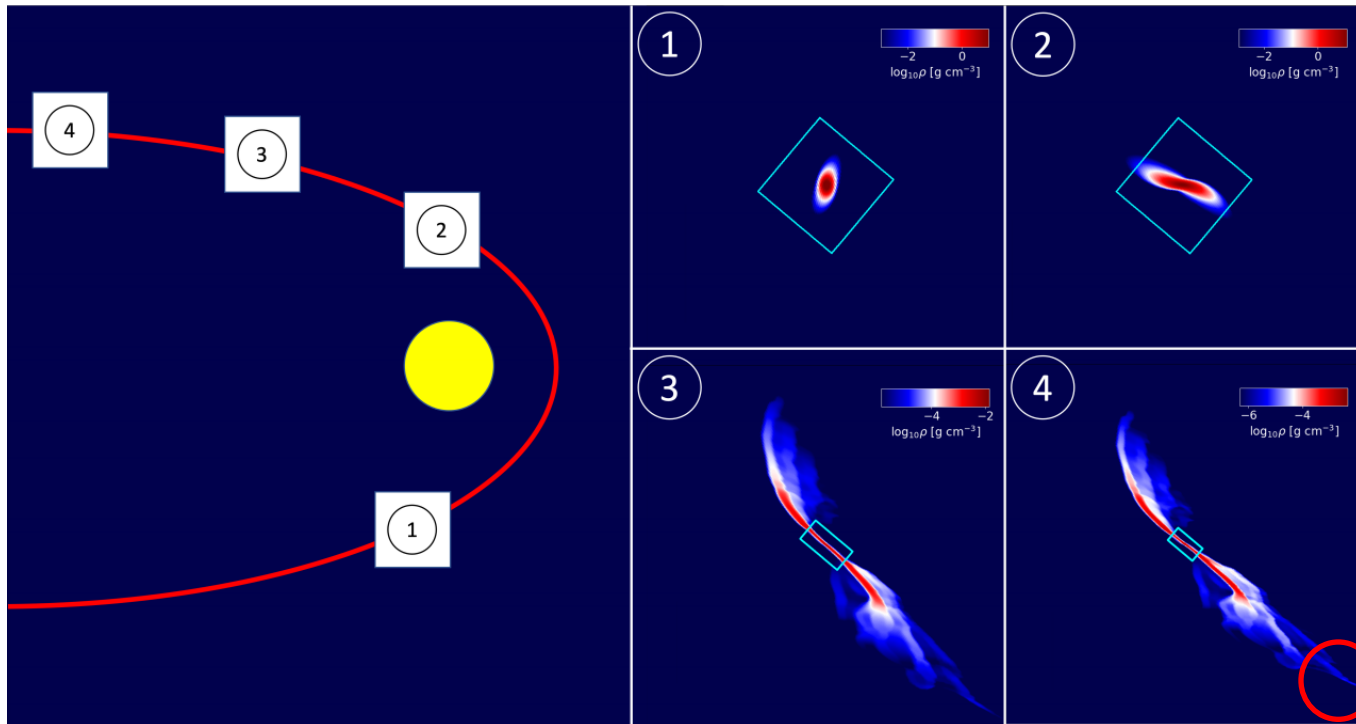
Milestones in energy (see the Conclusion section of this paper)

- Typical TDE flare radiates  $3 \times 10^{50}$  erg
- Circularizing bound debris into a disk (size  $\sim 2r_p$ ), shall release  $7.5 \times 10^{51}$  erg
- Accretion of all bound debris shall release  $3 \times 10^{53}$  erg

# Proposed Solutions to Inverse Energy Crisis

- within the framework of classic picture (Piran+2015)
  - Photon trapping in the accretion flow
  - Outflow (kinetic energy)
  - Outflow (blow the debris away)
  - Outflow (reprocess higher energy photon)
- Alternative possibility (Shiokawa+2015; Piran+2015; Krolik+2016)
  - Circularization is slow
  - Radiation is powered by shocks at the pericenter and apocenter

# Simulation setup



**Time unit** in the presentation

$$t_0 = \frac{\pi}{\sqrt{2}} \frac{GM_\bullet}{\Delta E^{3/2}} \simeq 7.6 \text{ days} \left( \frac{\Xi}{1.64} \right)^{-3/2} \left( \frac{M_\bullet}{10^5 M_\odot} \right)^{1/2} \left( \frac{M_\star}{3 M_\odot} \right)^{-1} \left( \frac{R_\star}{2.4 R_\odot} \right)^{3/2}$$

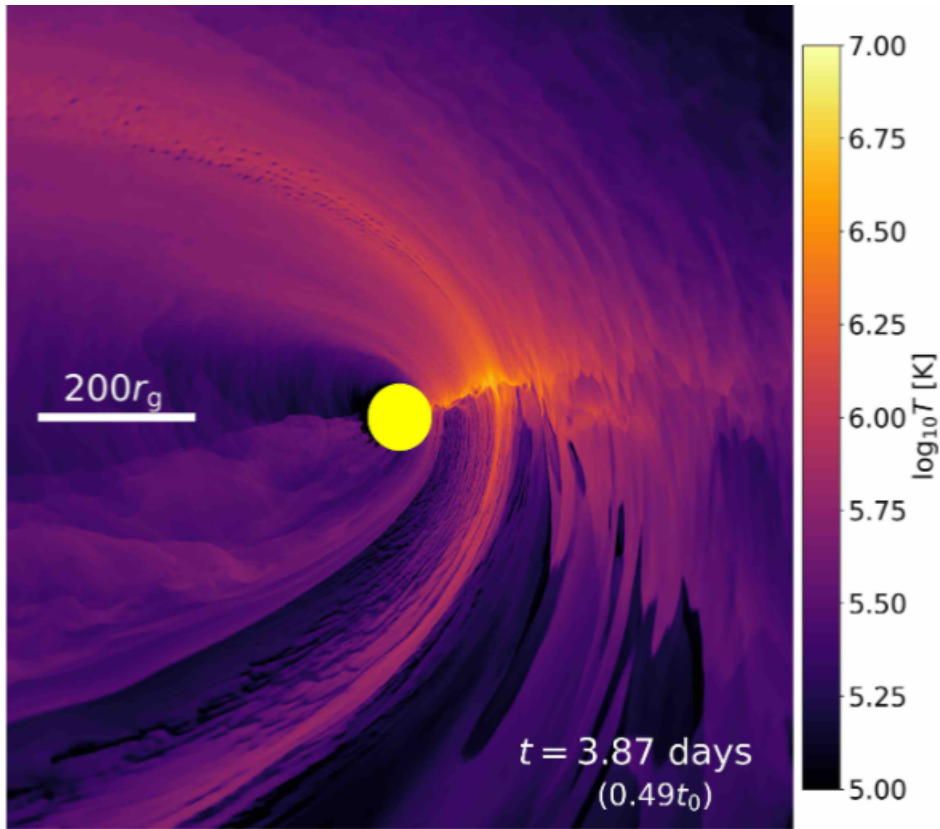
**Timing start point  $t = 0$ :** when the most tightly bound debris return to the orbital pericenter

← Most tightly bound debris

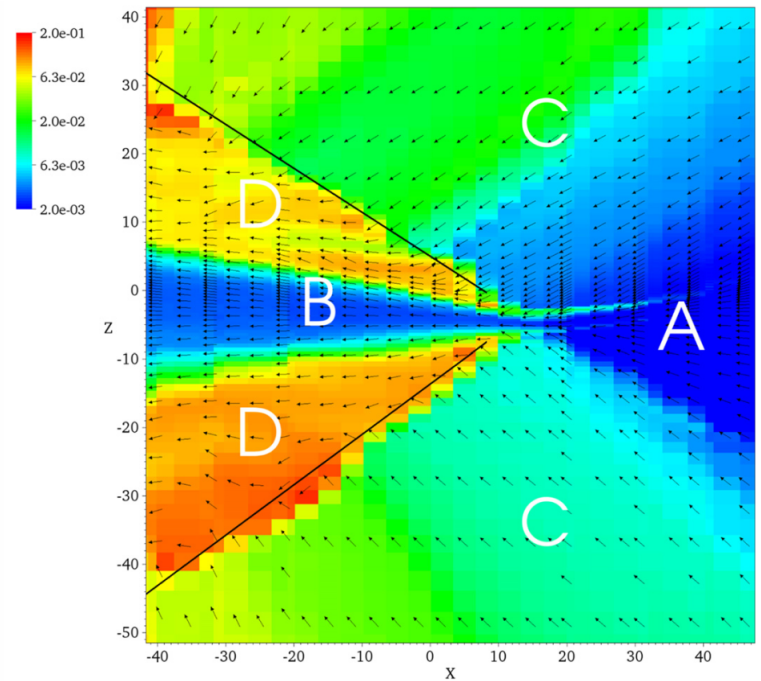
- fully relativistic hydrodynamic simulation
- 3 Msol star (MESA) +  $10^5$  Msol SMBH
- Long duration  $\sim 3$  weeks.
- Energy dissipated only through shocks, no viscosity

# The shock regions

**Nozzle shock:** debris stream converges to the equatorial plane



Temperature map in equatorial plane



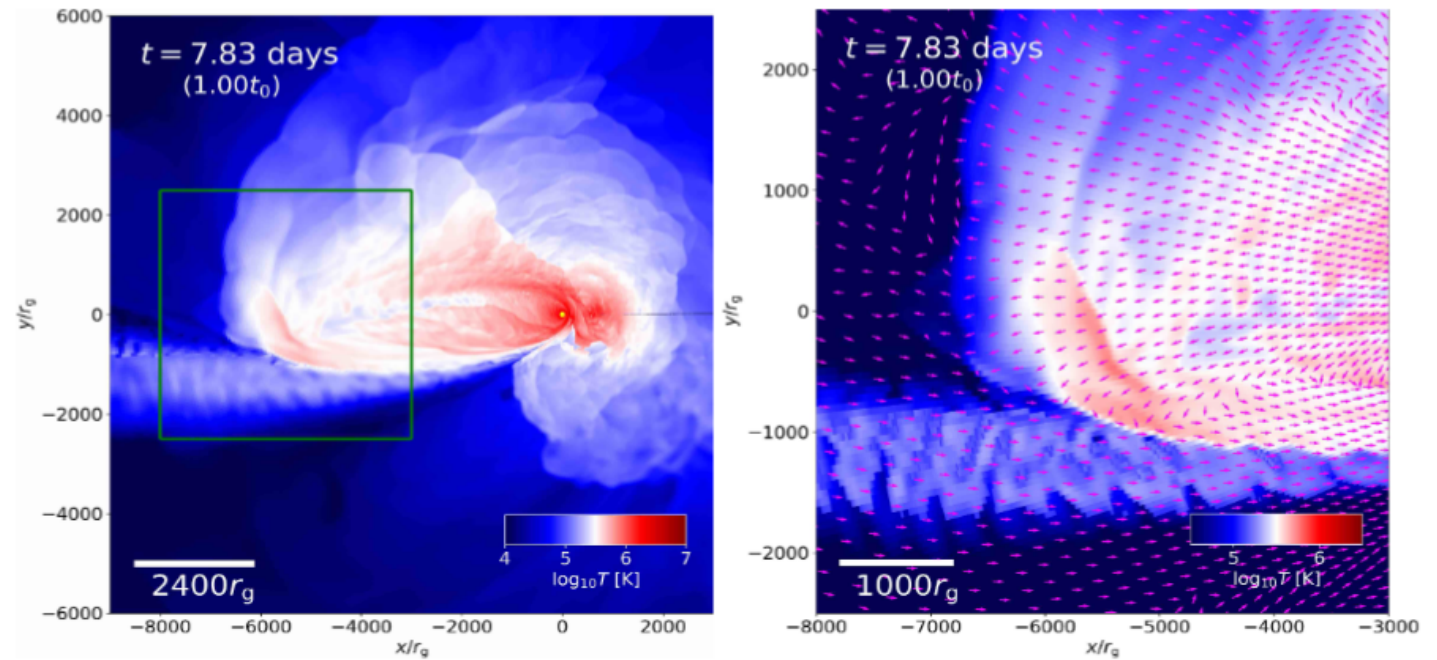
Entropy map in nozzle shock region (side view)



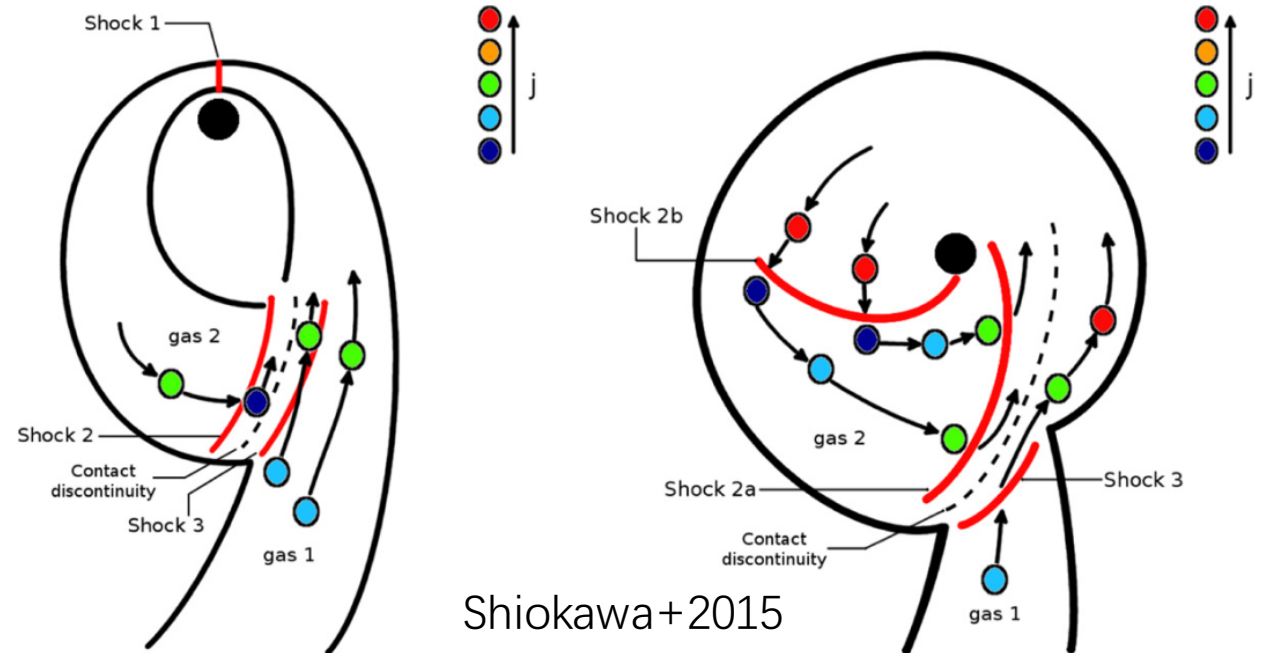
Shiokawa+2015

# The shock regions:

## Self-intersecting shock at apocenter



Note the shocks also redistribute the angular momentum of the fluid element, making the gas flow rounder and rounder.



Shiokawa+2015

# Circularized or not?

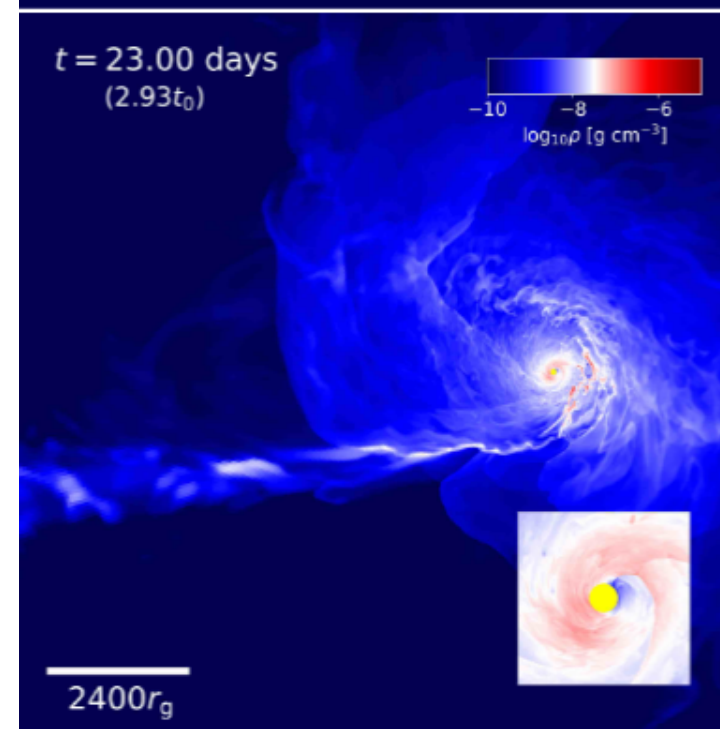
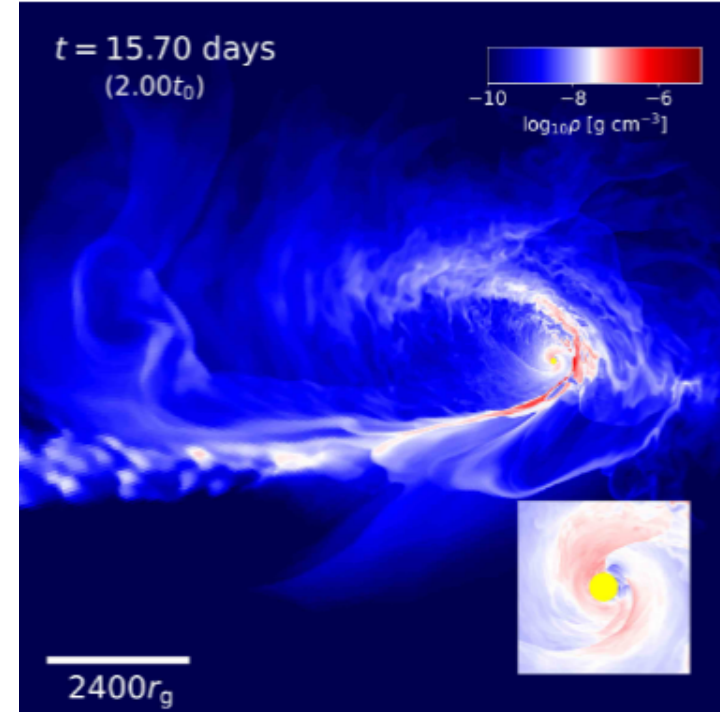
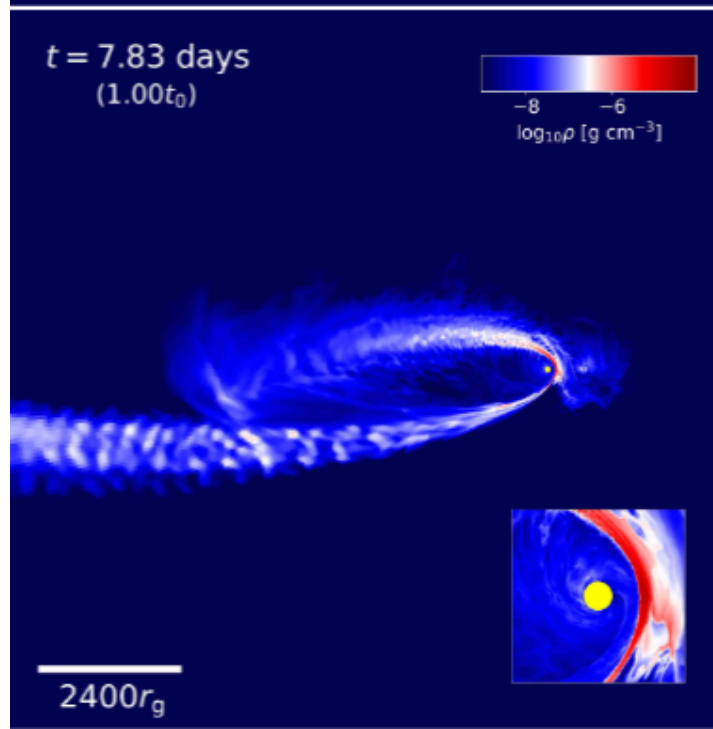
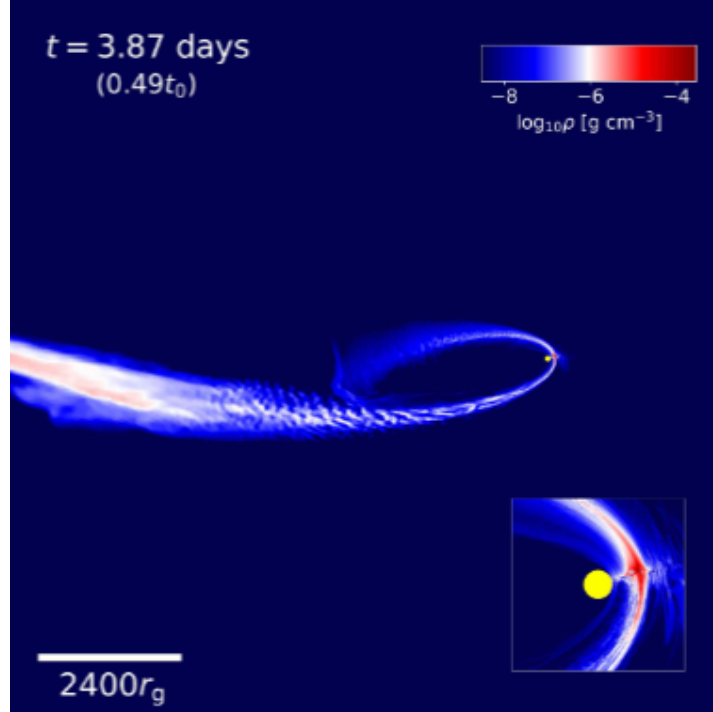
At the end of simulation ( $t \approx 3t_0$ ), the dissipated specific orbital energy is only **10%** of  $E_{\text{circ}}$

$$E_{\text{circ}} = \frac{GM_{\text{BH}}}{4r_p}$$

Just **enough to power the radiation during the simulated period.**

The debris forms an extended eccentric accretion flow with eccentricity  $\approx 0.4 - 0.5$

To fully circularize,  $> 30t_0$

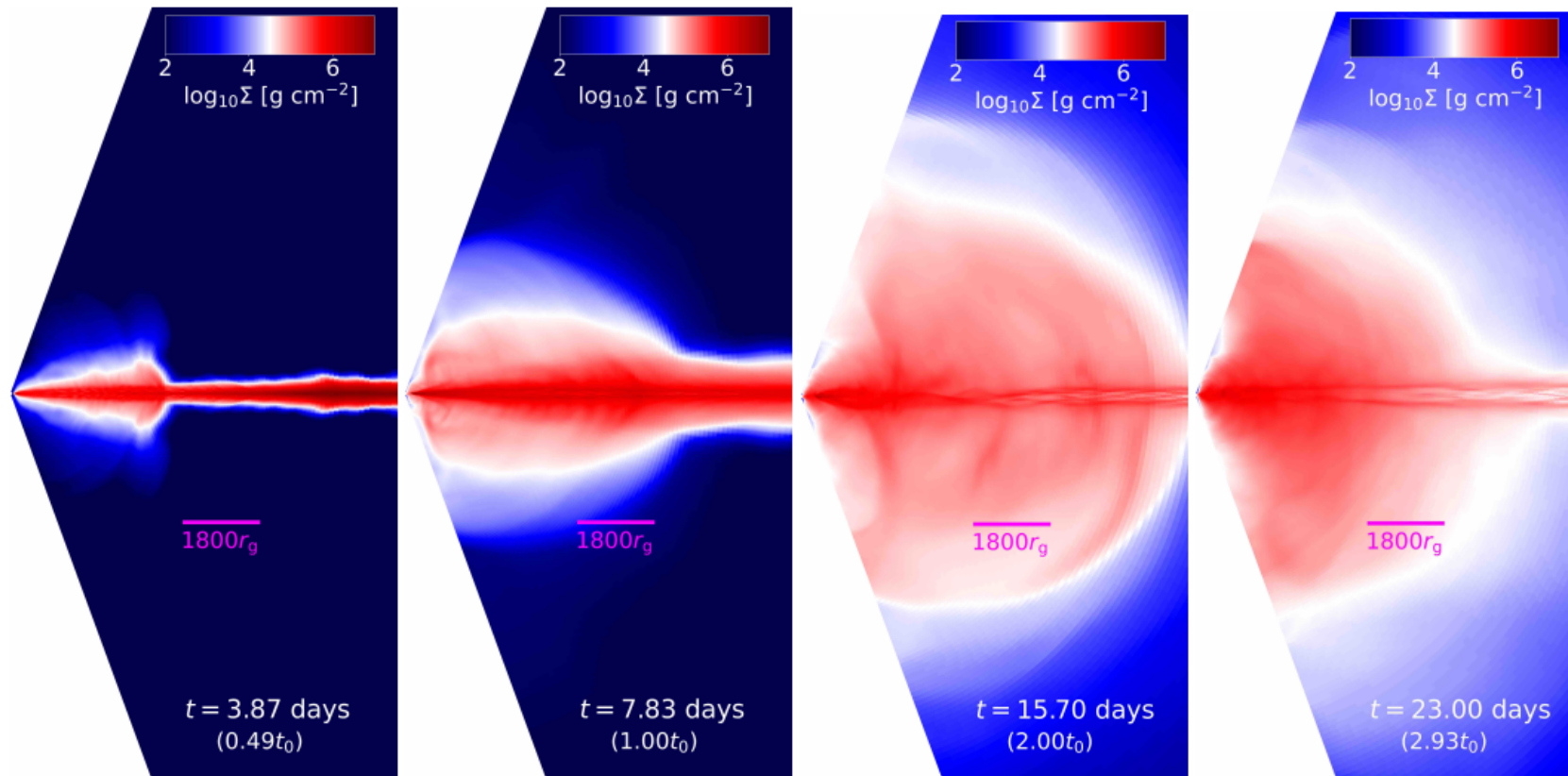




# Outflow?

The shocked gas expand outwards quasi-symmetrically, **marginally bound, and eventually falls back**

- Radiation pressure gradient built by shock heating
- Deflection caused by stream-stream collision



**Figure 6.** The azimuthally integrated density distribution at  $t/t_0 = 0.5, 1, 2,$  and  $3.$

# Photosphere radius

Thermalization photosphere

$$\sqrt{\tau_T \tau_{ff}} \simeq 1$$

Generally quasi-spherical in shape, but also depend on the azimuthal angle  $\Phi$

At  $t = t_0$ ,  $r \simeq 4000 - 5000 r_g$ .

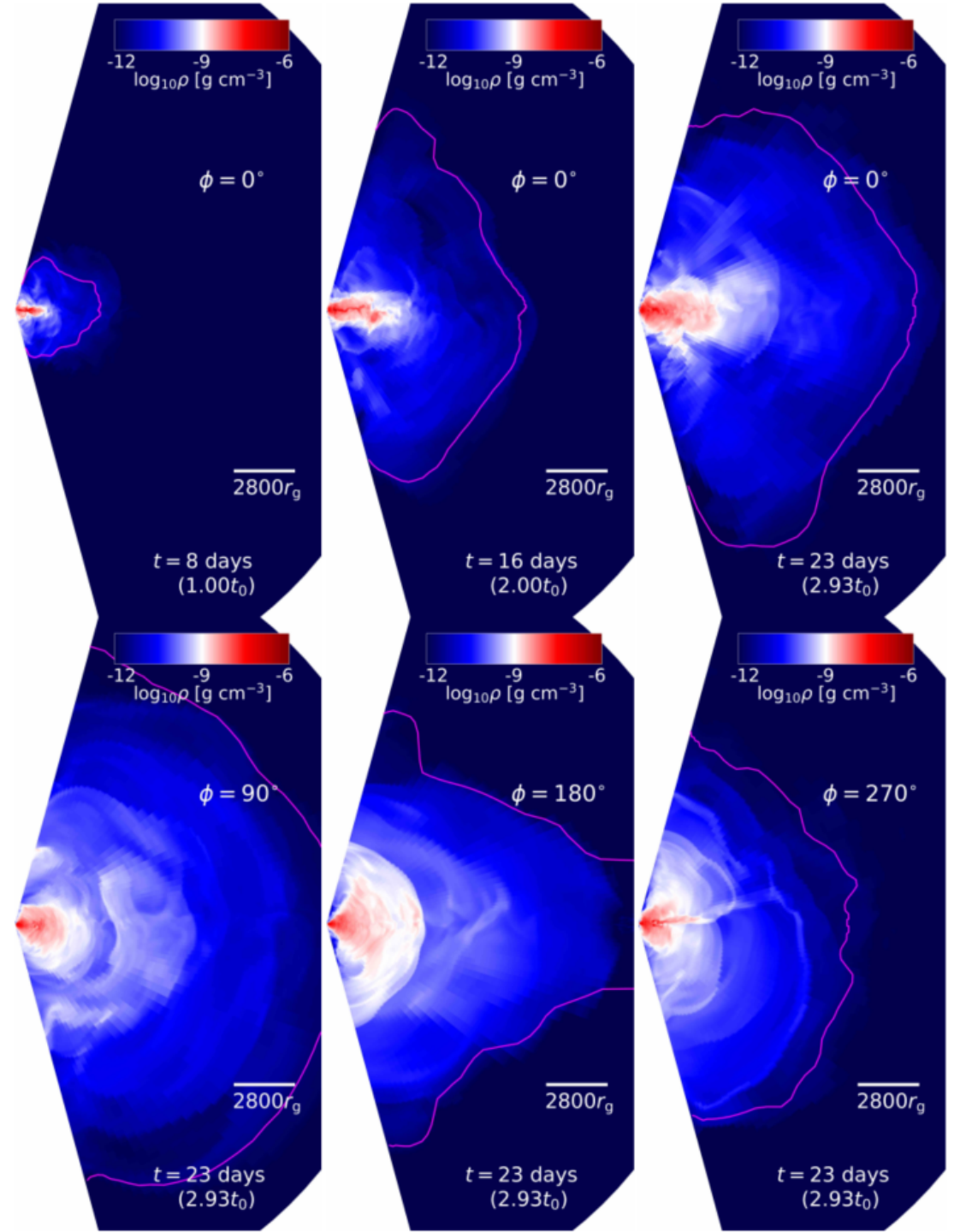
$t = 2t_0$   $9000 - 10000 r_g$

$t = 3t_0$   $\simeq 12000 r_g$

$$r_g = 1.48 \times 10^{10} \text{ cm}$$

$$L = \int_0^{2\pi} \int_{\theta_c}^{\pi - \theta_c} \int_{r=r(t_{\text{cool}} < t)}^{r=r(\tau=1)} \frac{aT^4}{t_{\text{cool}}} r^2 \sin \theta dr d\theta d\phi,$$

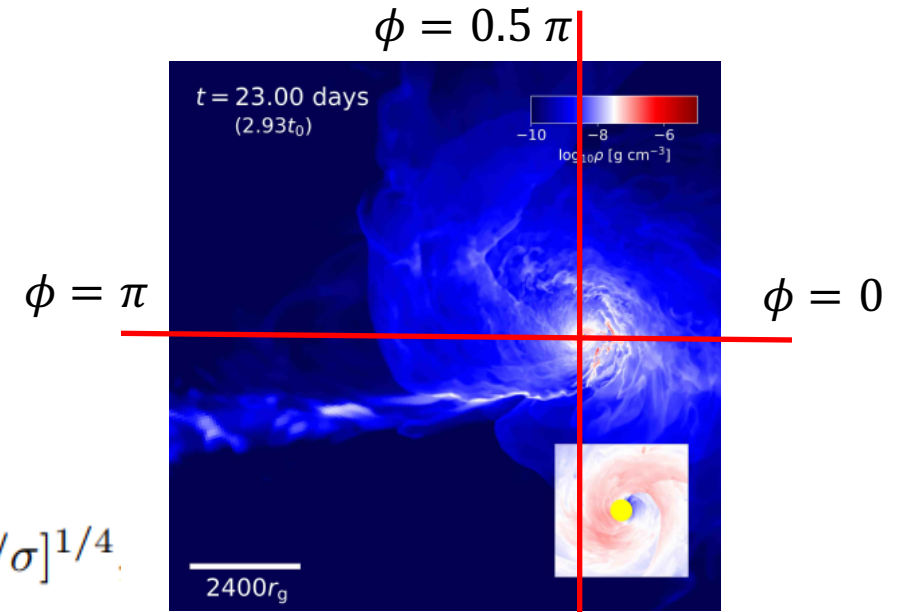
We estimate that the peak luminosity is  $\simeq 10^{44}$  erg/s  $\simeq 10 L_{\text{Edd}}$ , which occurs at  $t \simeq t_0$ . This is roughly the mean rate of thermal energy creation during the simulation. The photospheric temperature distribution at



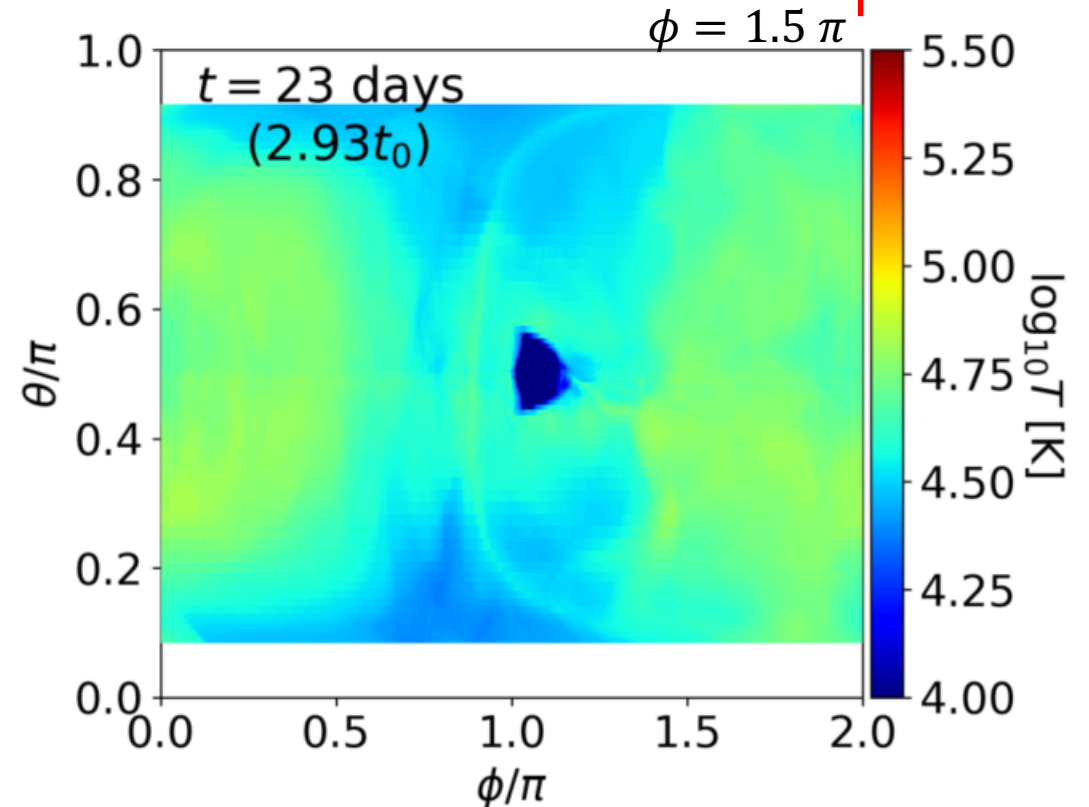
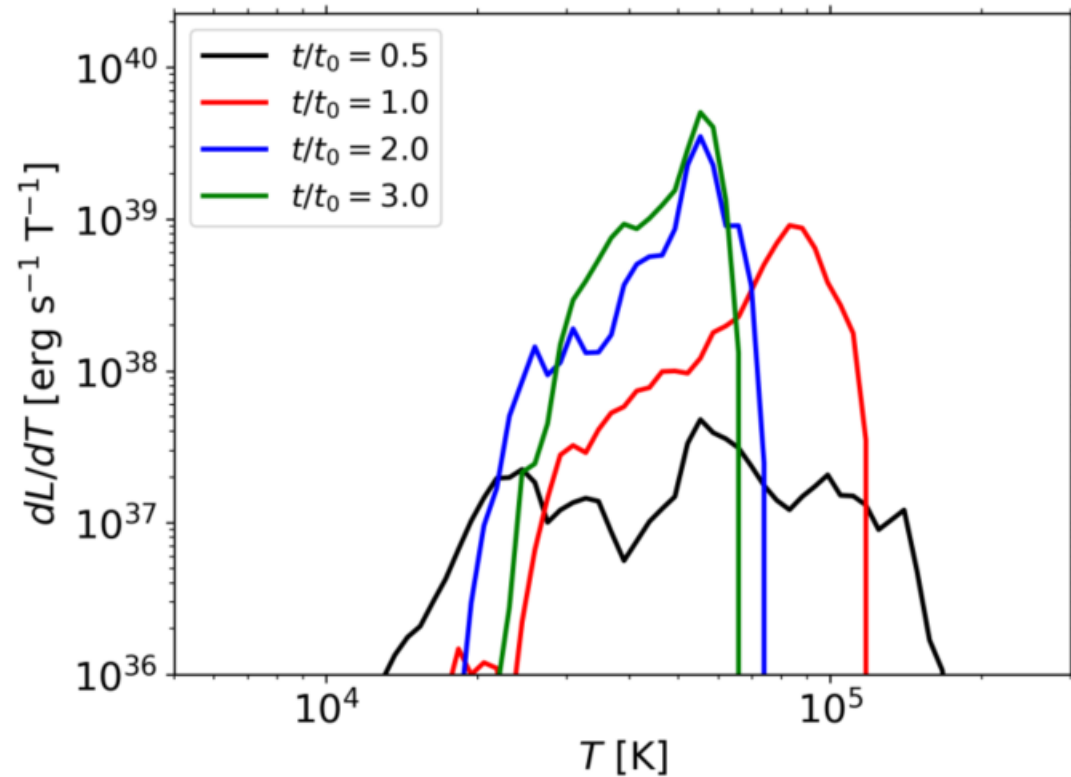
# Temperature

Interesting points:

- Multi-temperature photosphere
- Observed temperature depends on view angle?



$$T_{\text{ph}} = [(dL/dA)/\sigma]^{1/4}$$



# Conclusions

1. Shocks power the TDE radiation (at least in the simulated period)
2. Swift “circularization” does not happen, (need at least  $30t_0$ )
3. Radiatively efficient accretion of most of the debris mass onto the black hole certainly did not happen.
4. The initially bound debris does not become unbound.