

**Wind-fed GRMHD simulations of Sagittarius
A*: tilt and alignment of jets and accretion
discs, electron thermodynamics, multiscale
modelling of the rotation measure**


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(Several slides taken from talk given by Ressler at ITC, Harvard University)

December 20, 2023

Image Credit: Event Horizon Telescope Collaboration 2022

Sagittarius A* at 230 GHz

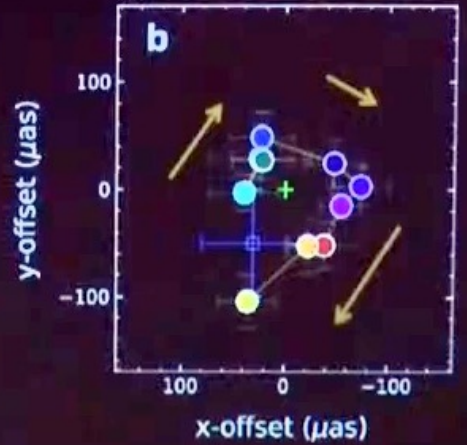
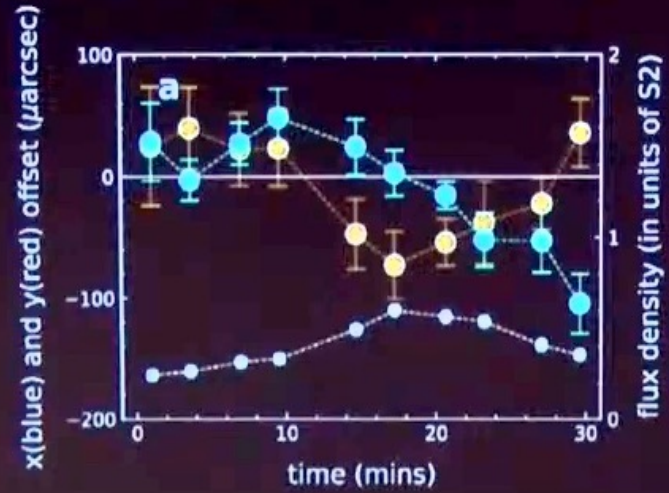
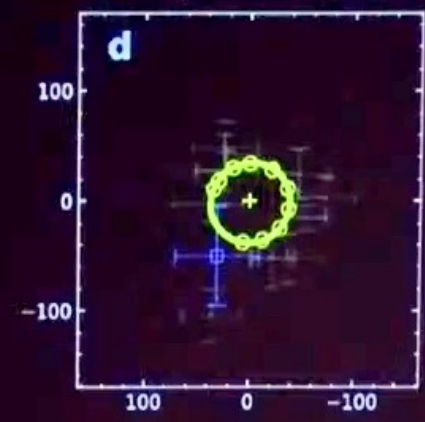
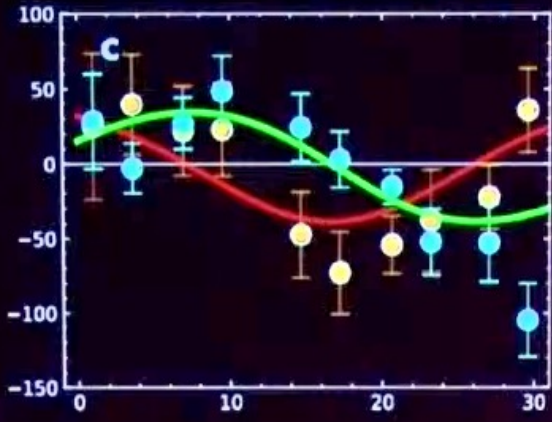


GRAVITY (NIR)

Jul 22 2018 flare, MJD=58321.9954

$R=7 R_g$ $a=0$ $i=160^\circ$ $\Omega=160^\circ$ $\chi_r^2=1.2$

Flare motion
at $\sim 5 r_s$



GRAVITY
Collaboration 2018

Horizon Scale Simulations

Conservation of Mass

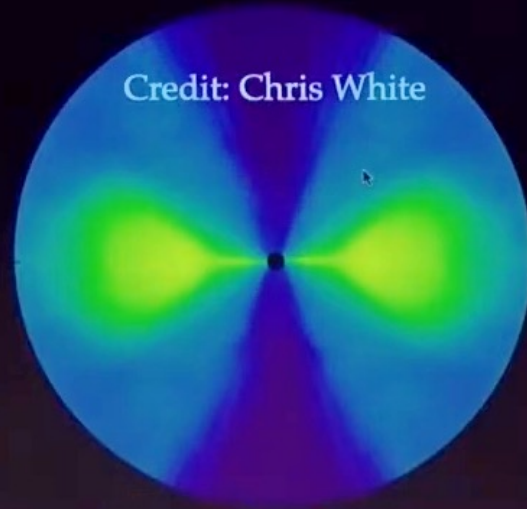
$$\nabla_{\mu}(\rho u^{\mu}) = 0$$

Conservation of Energy/Momentum

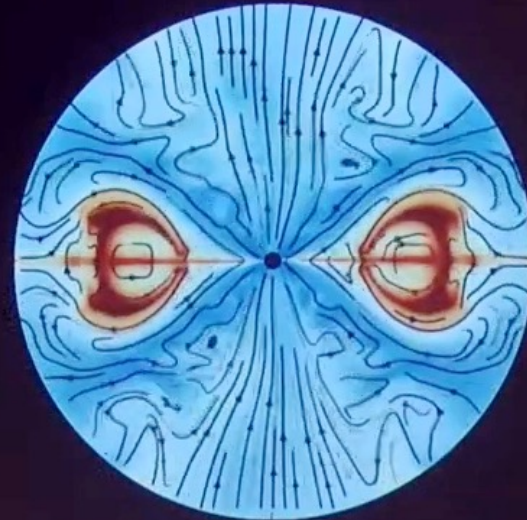
$$\nabla_{\mu} T^{\mu}_{\nu} = 0$$

Maxwell's Equations

$$\partial F^{*\mu\nu} = 0$$

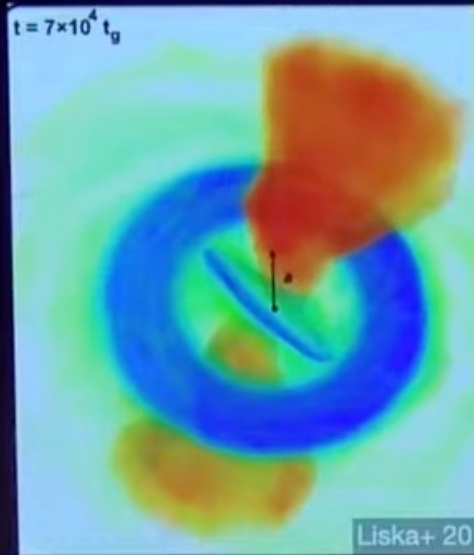


Mass Density



$P_{\text{mag}} / P_{\text{gas}}$

Horizon Scale Simulations



What we know:

$M \sim 4 \text{ Million } M_{\odot}$

What we are pretty sure about:

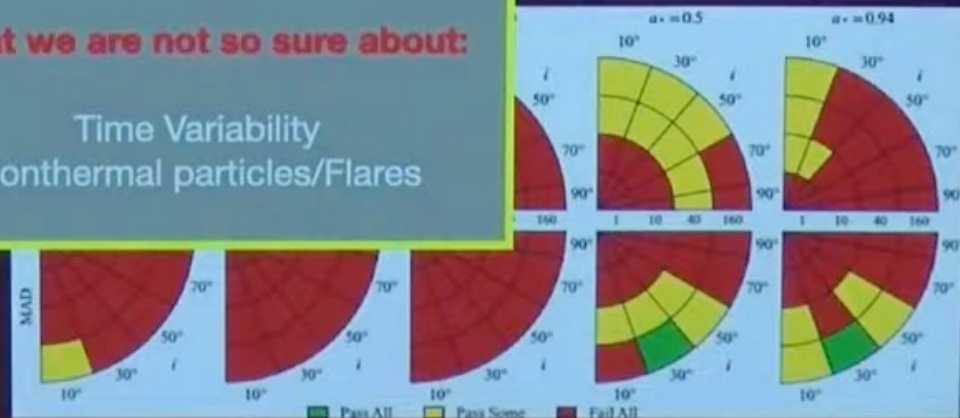
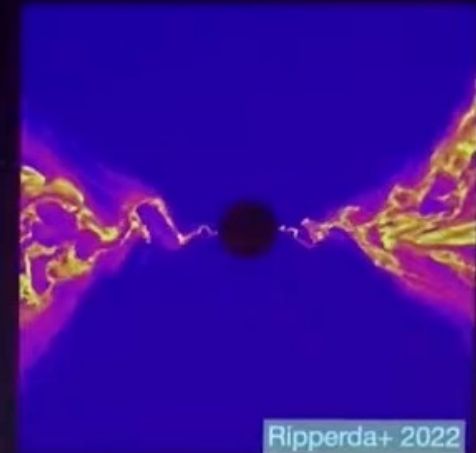
Accretion flow \sim face on near horizon
Strong B-field near horizon

Models Favor:

Higher black hole spin
Magnetically Arrested

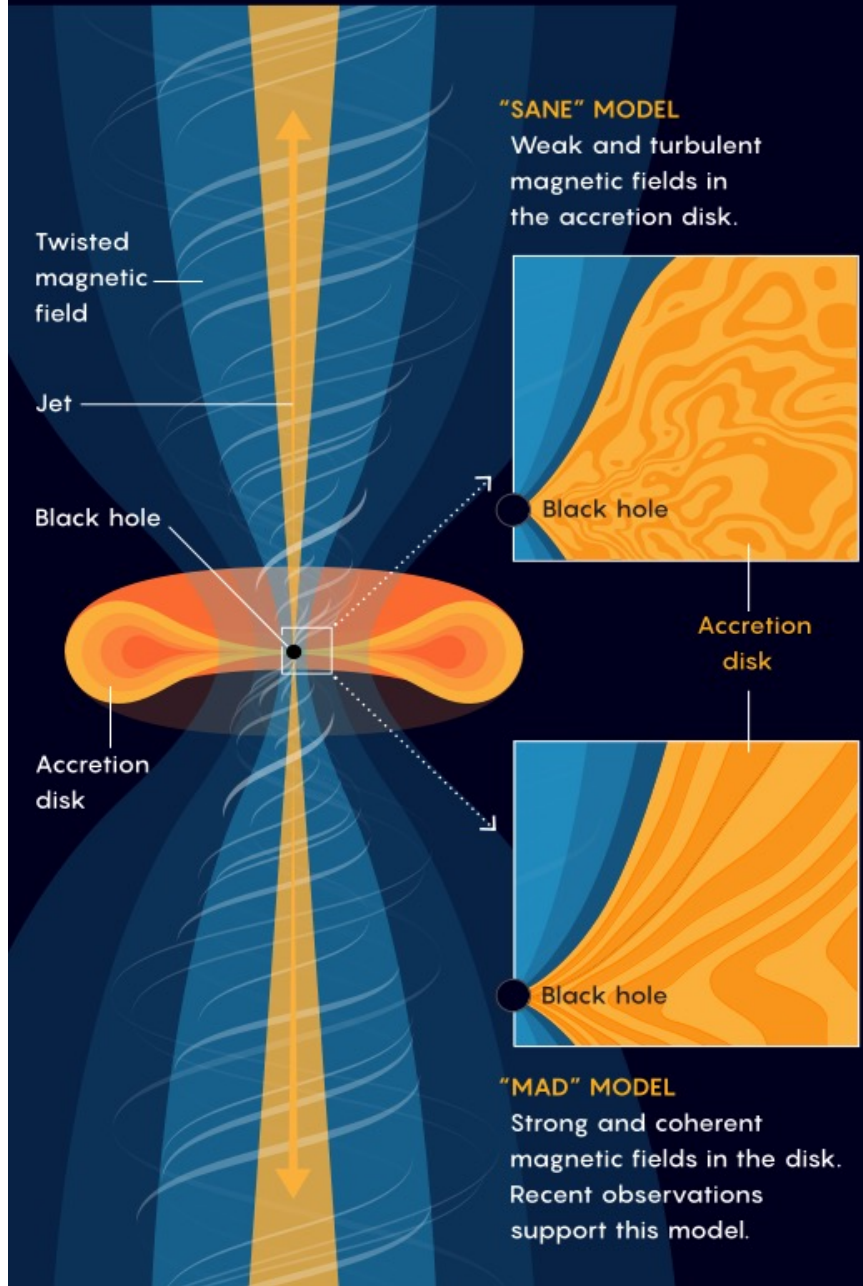
What we are not so sure about:

Time Variability
Nonthermal particles/Flares



Inside a Black Hole's Jet Engine

As a spinning black hole pulls in matter, it creates a rotating "accretion disk" of charged particles. The motion generates twisted magnetic fields that accelerate particles into two thin jets.



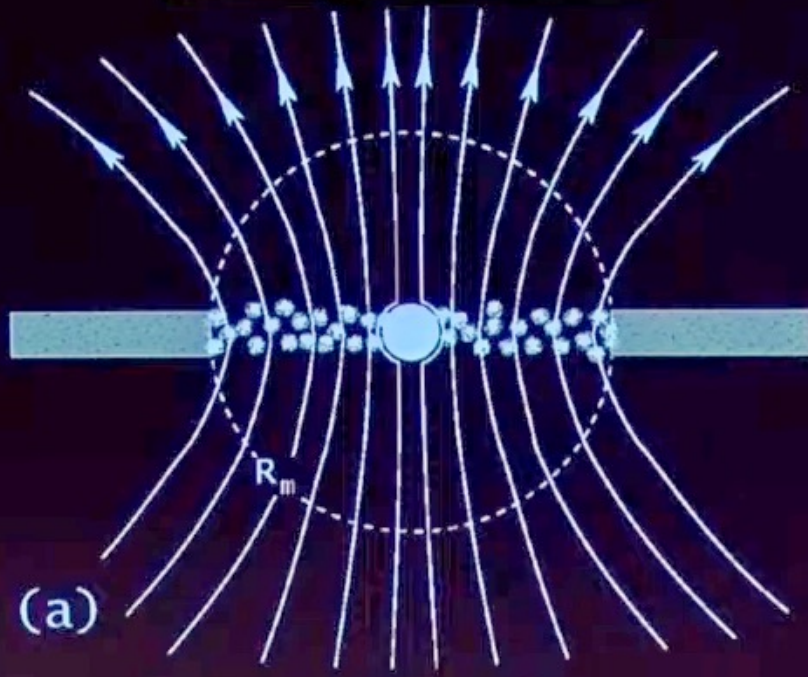
Standard And Normal Evolution (SANE)

Magnetically Arrested Disk (MAD)

Image credit: Quanta Magazine

Magnetically Arrested Disks (MADs)

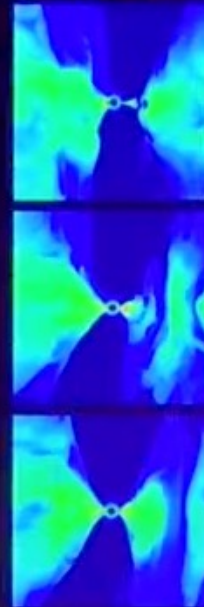
Flux Accumulation



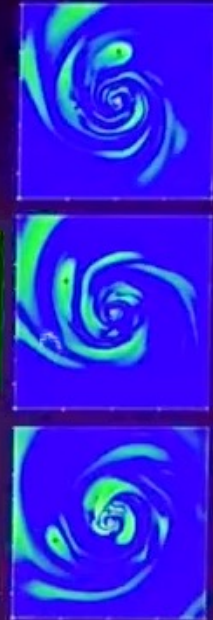
(a)

Narayan+ 2003

Chatterjee & Narayan 2022



Porth+ 2020



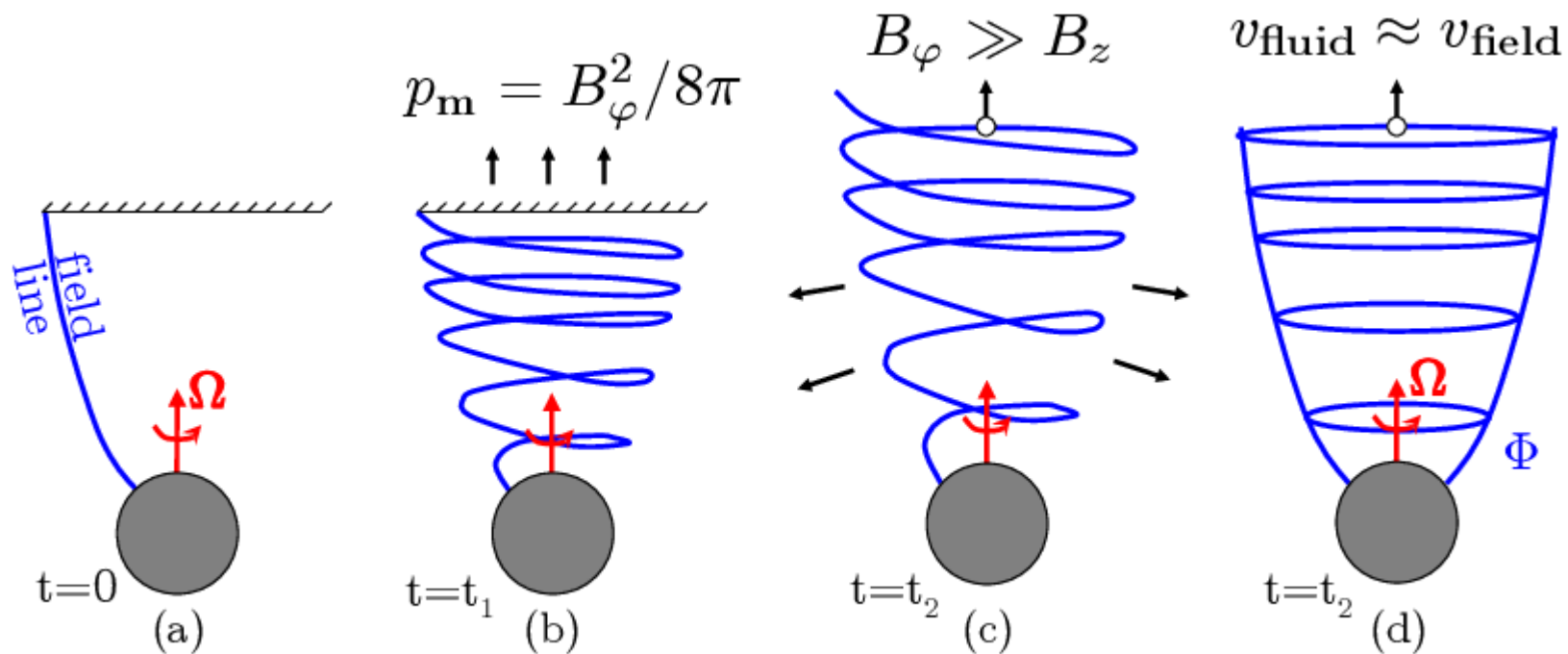
Flux Eruption

$$\frac{\Phi_{\text{BH}}}{\sqrt{\dot{M} r_g^2 c}} \approx 50-70$$



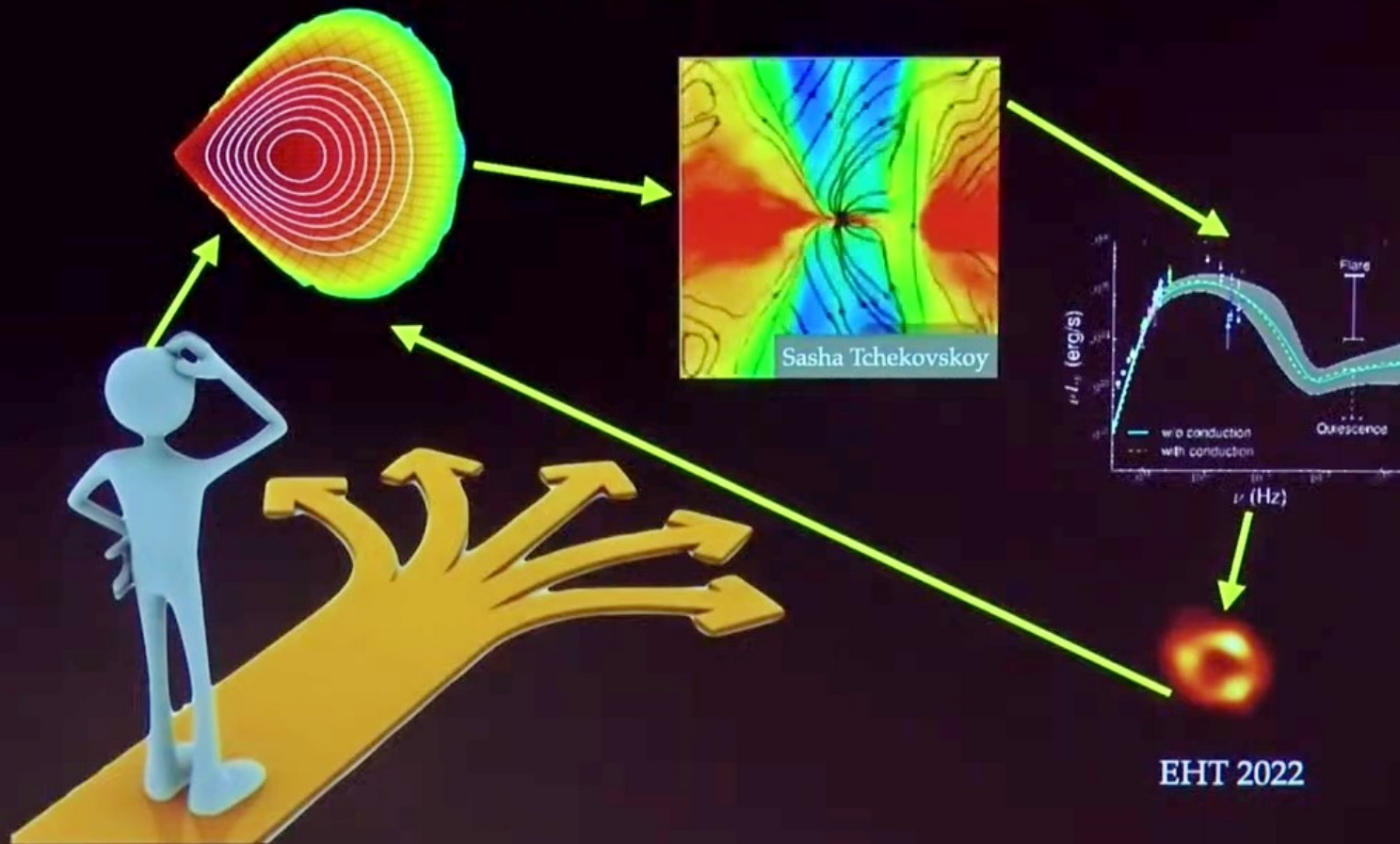
Tchekovskoy+ 2011

Blandford-Znajek (BZ) process for spinning black hole
 (Davis & Tchekhovskoy, ARAA, 2020)



4. Illustration of jet formation by magnetic fields (a) Consider a purely vertical

Constraining Accretion Physics From Observations



Bondi Radius Resolved

Hot gas from
Wolf-Rayet
Stars and
other point
sources

$T \sim 1.3 \text{ keV}$
 $n_e \sim 26 \text{ cm}^{-3}$

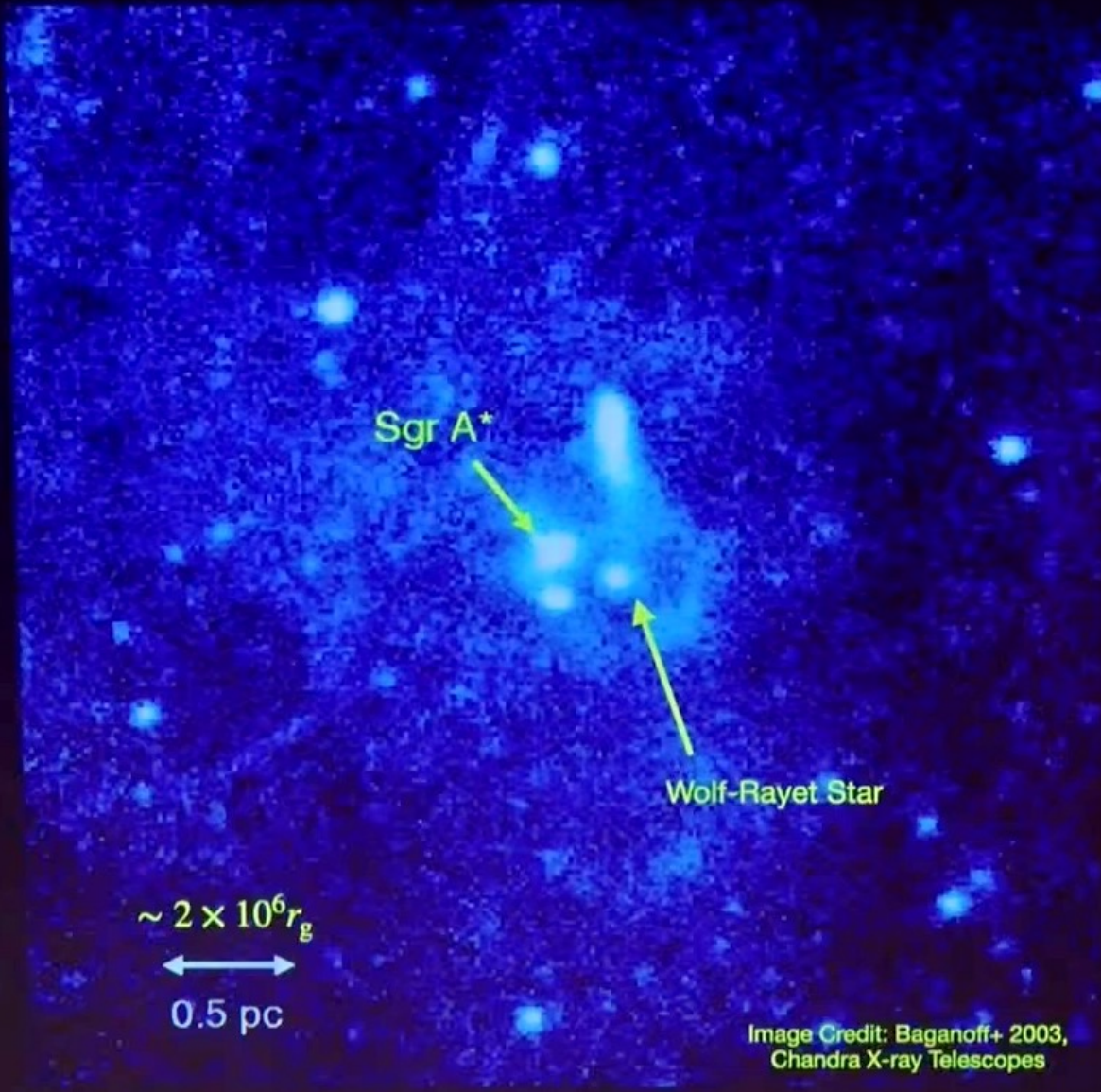
$\dot{M}_{\text{Bondi}} \sim 10^{-6} M_{\odot}/\text{yr}$

$\sim 2 \times 10^6 r_g$
↔
0.5 pc

Image Credit: Baganoff+ 2003,
Chandra X-ray Telescopes

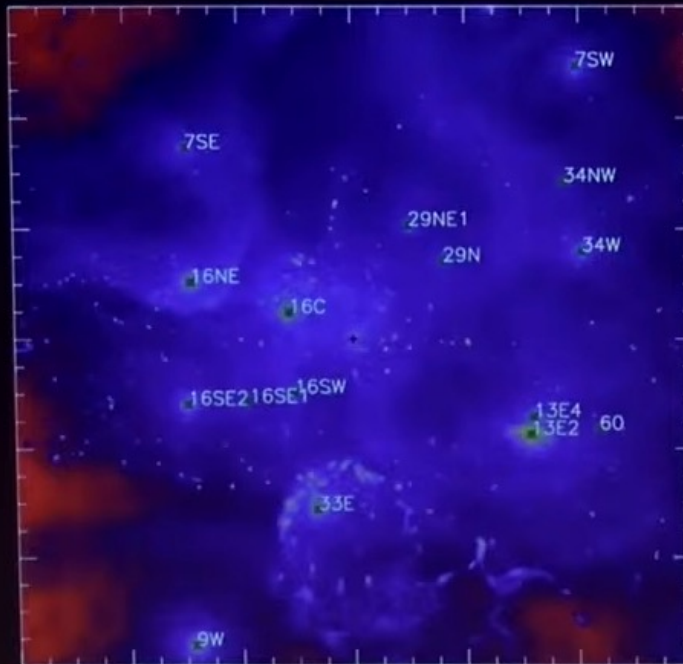
Sgr A*

Wolf-Rayet Star

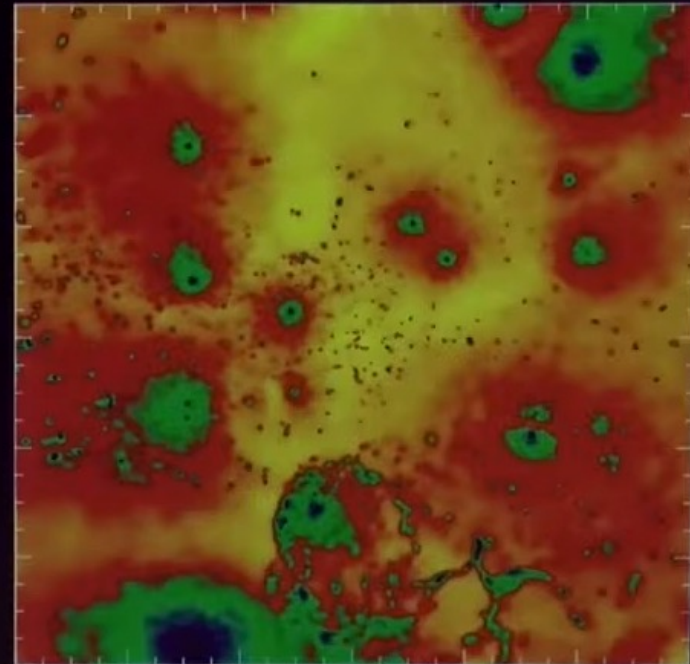


Bondi Scale Simulations

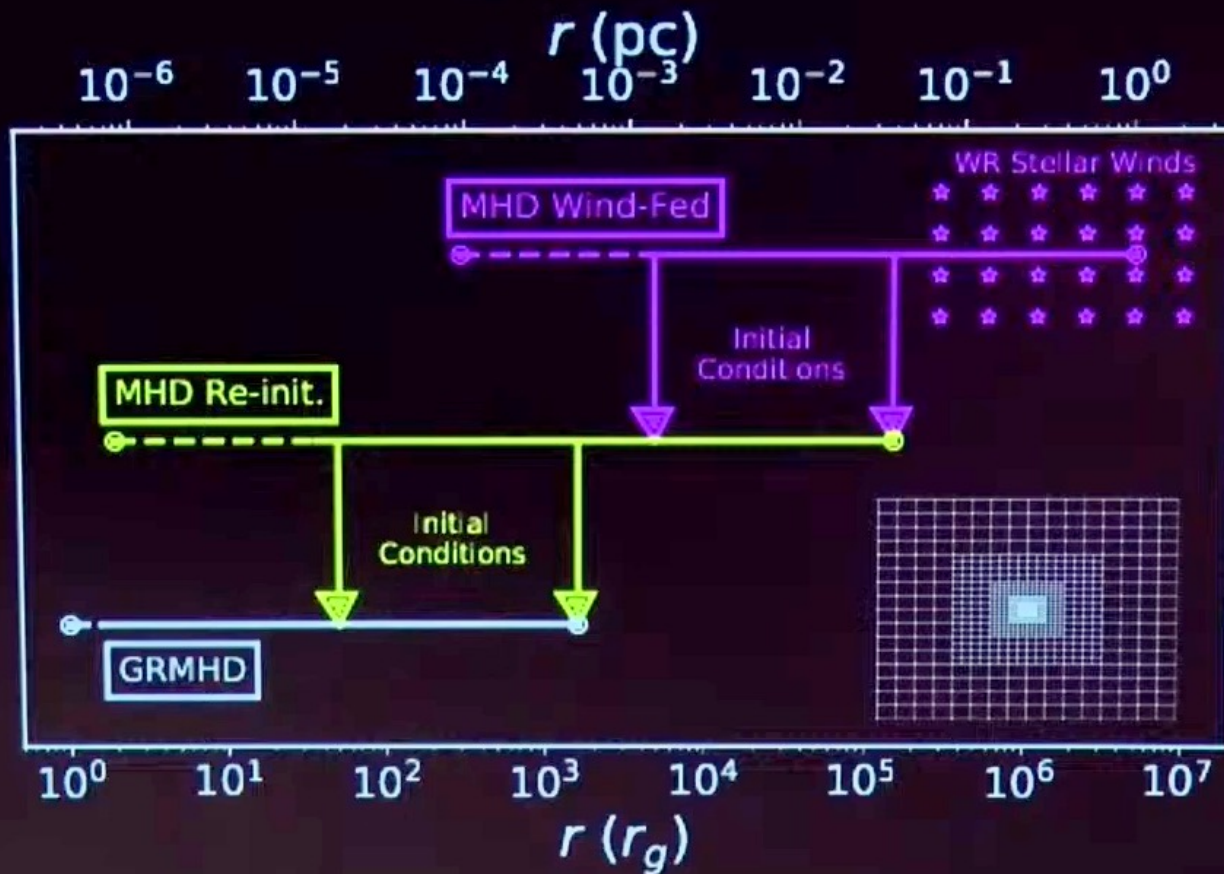
Cuadra+ 2005, 2006, 2008
Calderon+ 2019
Russell+ 2019




~ 0.2 pc



Bridging the Gap



Main Parameters of the Model

$$\beta_w \equiv \frac{P_{\text{ram}}}{P_{\text{mag}}}$$




t_{restart}



$$f_e = f_e(\beta, T_e/T_i, \sigma, \dots)$$



Electron Thermodynamics

$$\nabla_{\mu}(\rho u^{\mu}) = 0$$

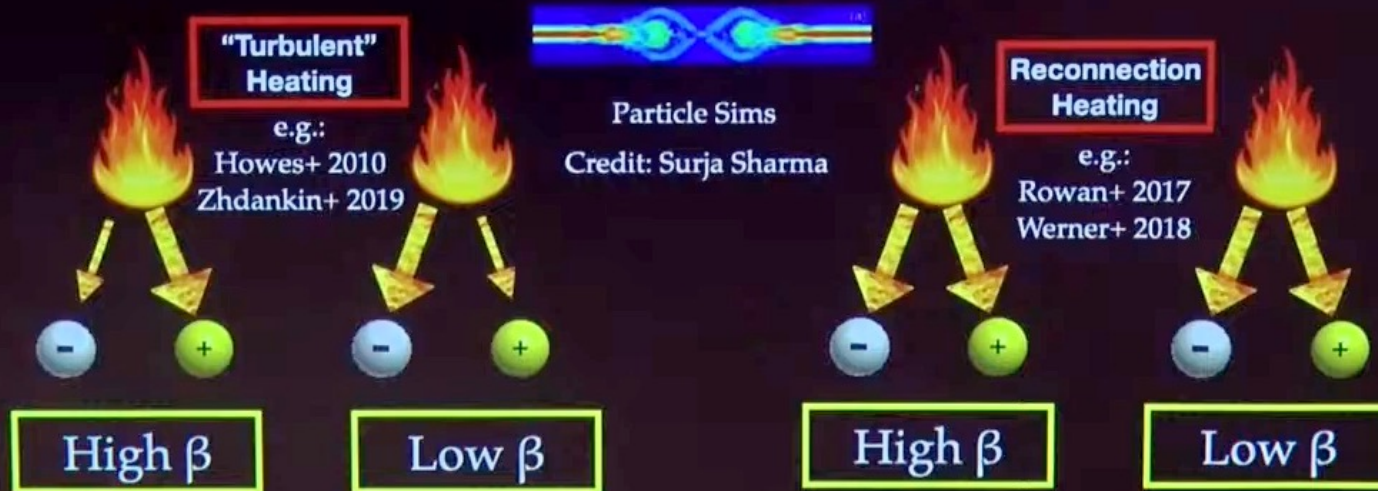
$$\nabla_{\mu} T_{\nu}^{\mu} = 0$$

$$\partial F^{*\mu\nu} = 0$$



Sgr A*
 $10^5 r_g \sim$ collisional
 $r_g \sim$ very collisionless

$$f_e(\beta, T_e, T_p, \dots)$$



Smorgasbord of Models

$$\beta_w = 10^2, a = 0$$

3 different realizations

$$\beta_w = 10^2, a = 0.9375$$

2 different realizations

$$\beta_w = 10^6, a = 0$$

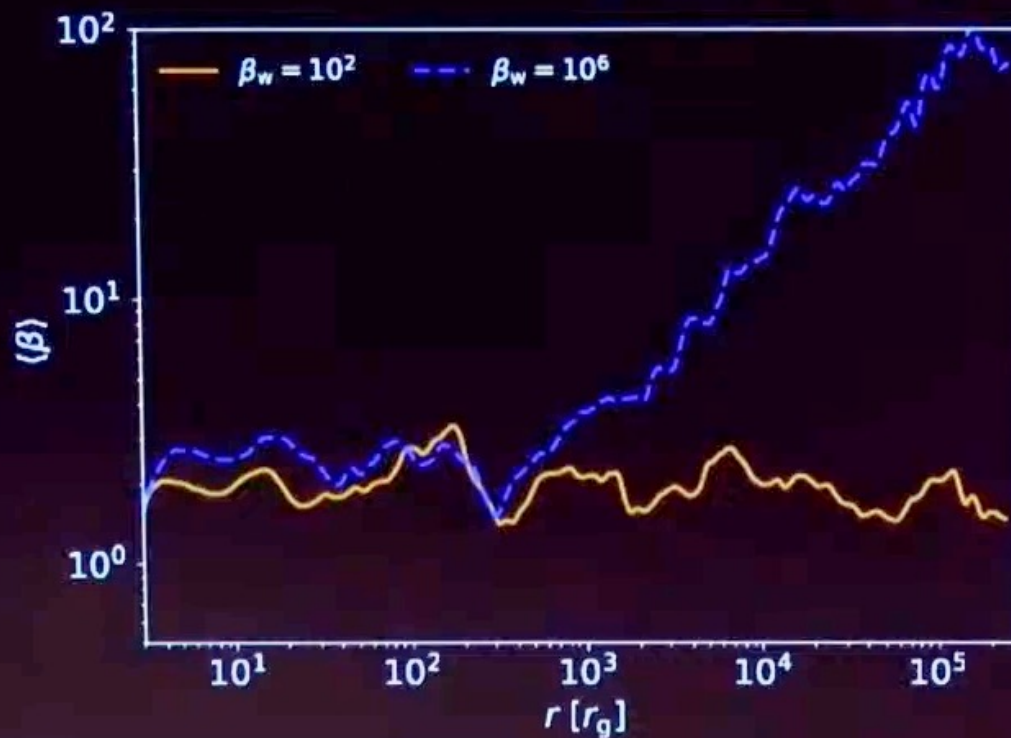
2 different realizations

$$\beta_w = 10^6, a = 0.9375$$

1 realization



Strongly Magnetized Flows Seem Inevitable

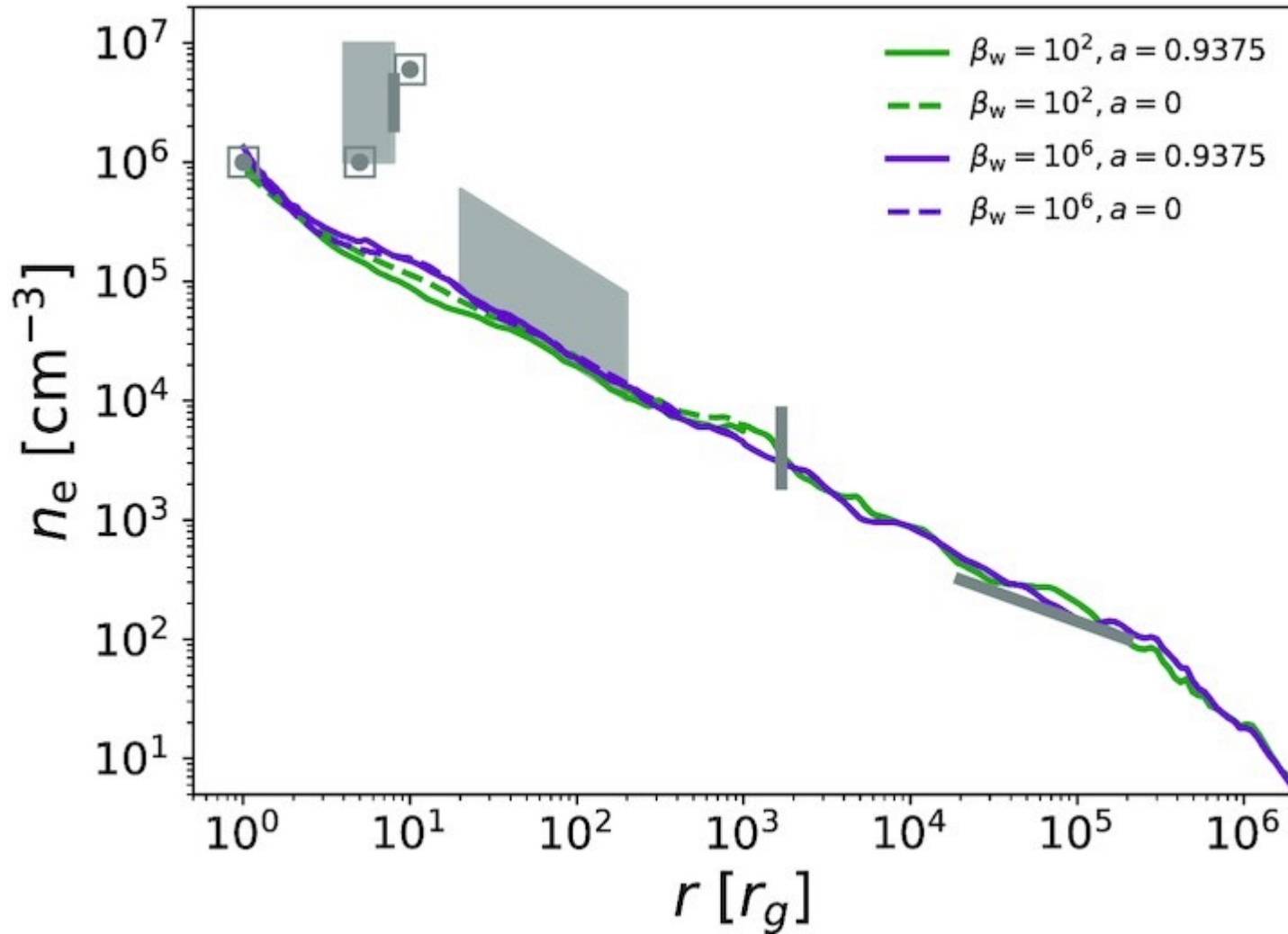


Weakly Magnetized Winds

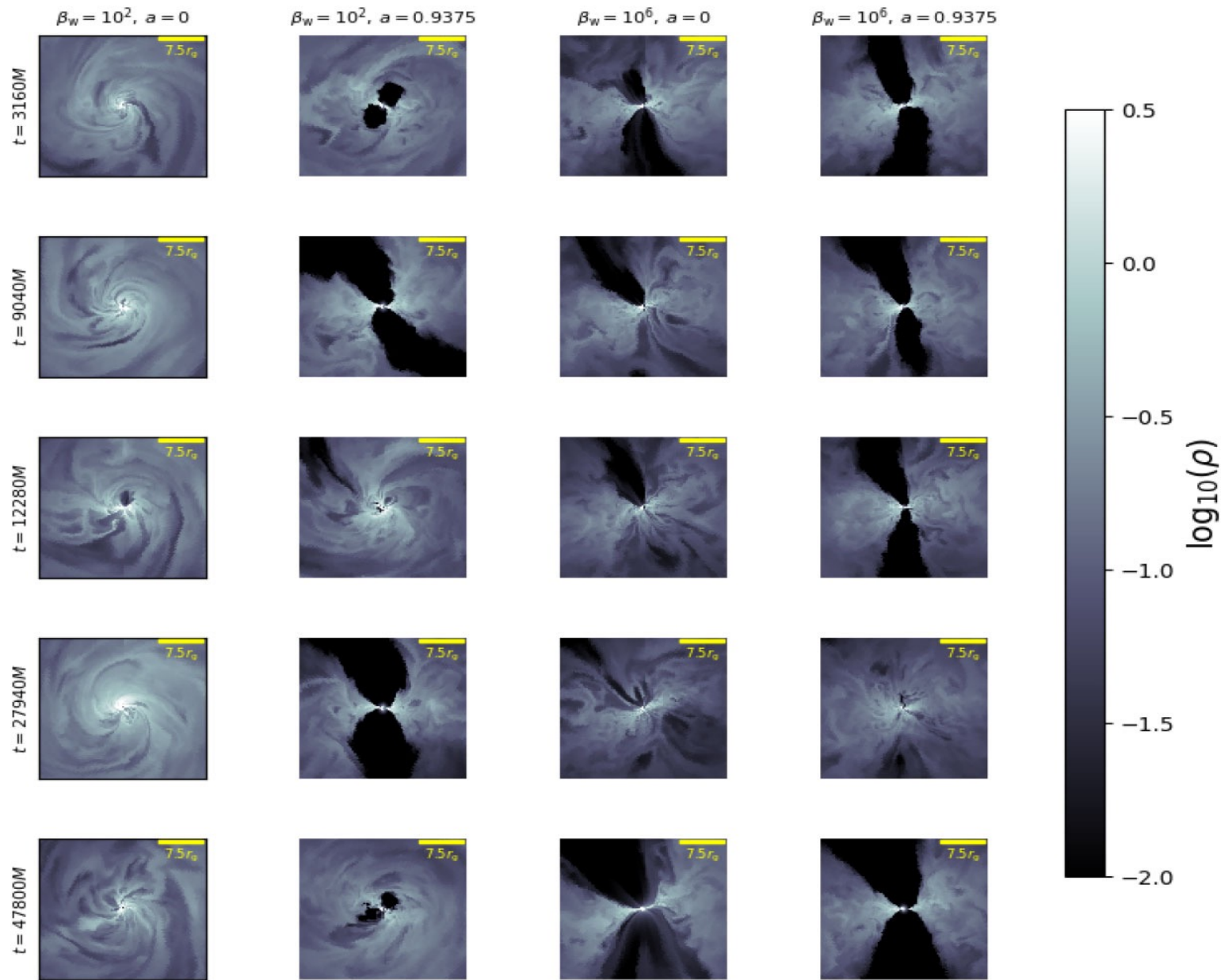
$$\beta \equiv \frac{P_{\text{g}}}{P_{\text{mag}}}$$

Strongly Magnetized Winds

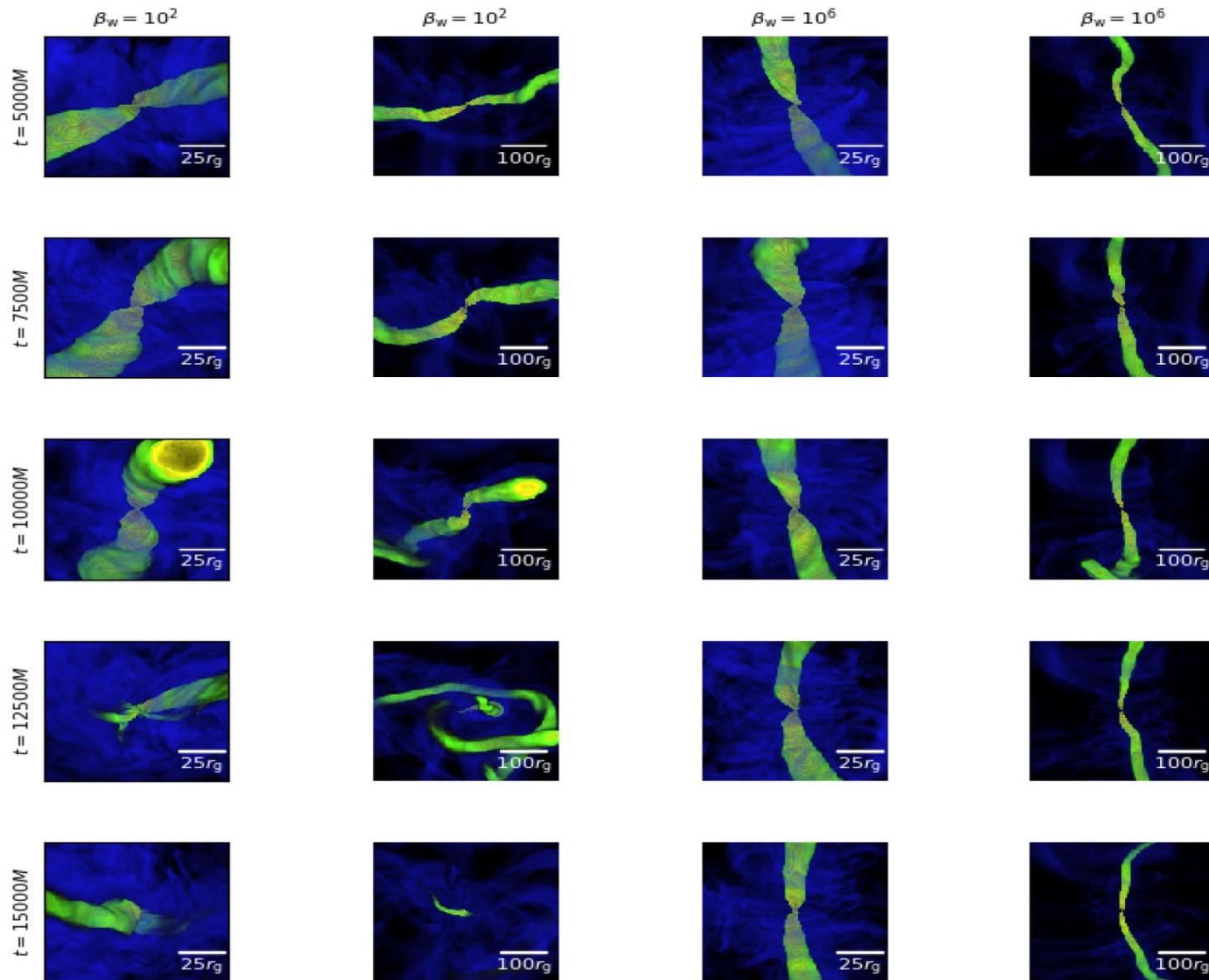
Density profile/accretion rate consistent with observations



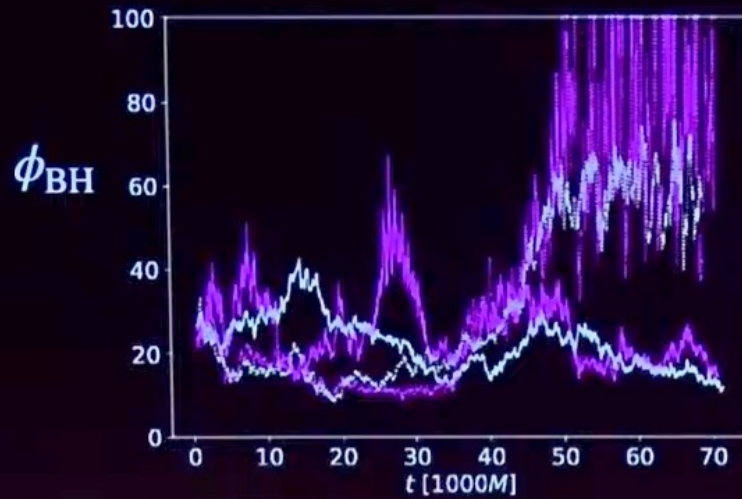
Mass Density in x-z plane



3D volume renderings of the jets with $a = 0.9375$ using magnetization parameter at five different times and two different spatial scales

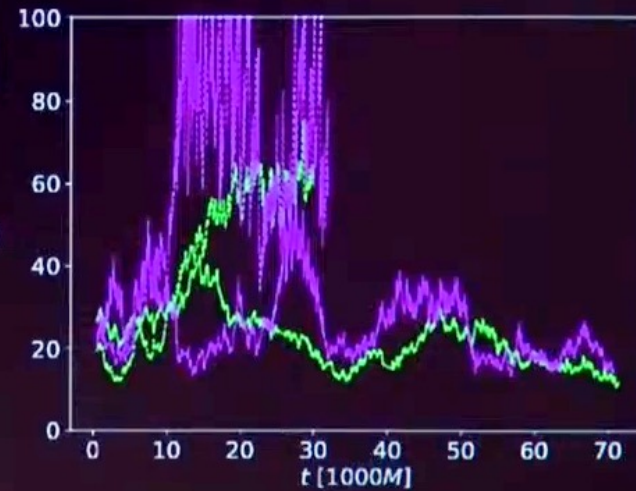


MAD Sometimes, Not Always



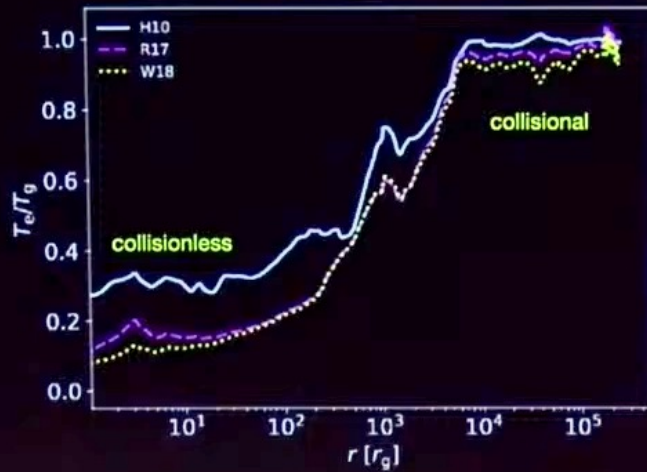
6/8 Simulations
Magnetically Arrested

ϕ_{BH}



Electron Temperatures

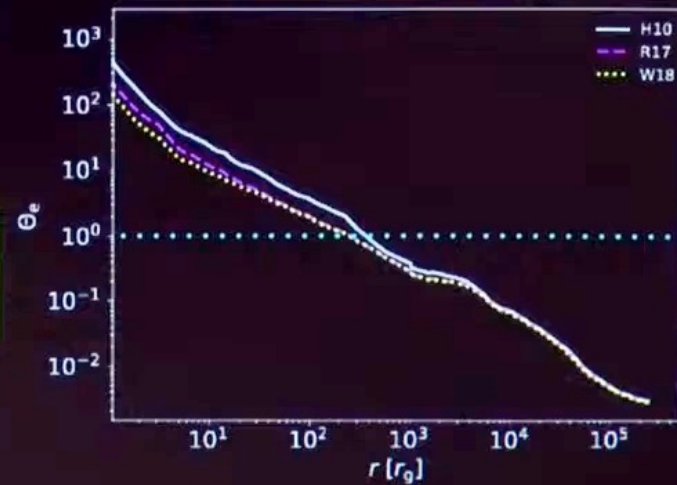
Electron Temperature Ratio



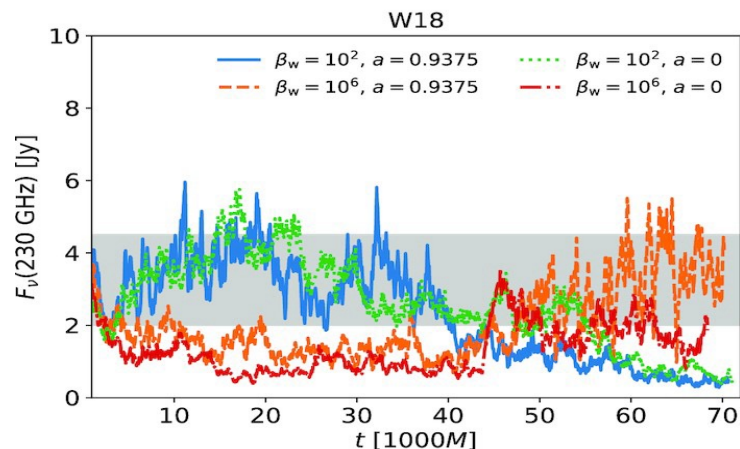
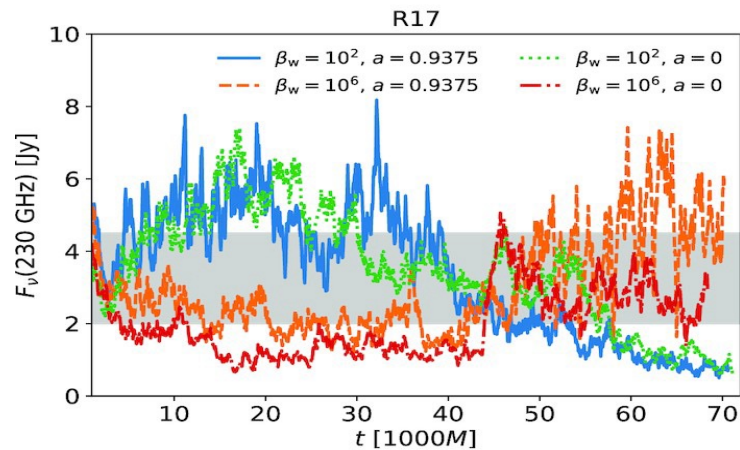
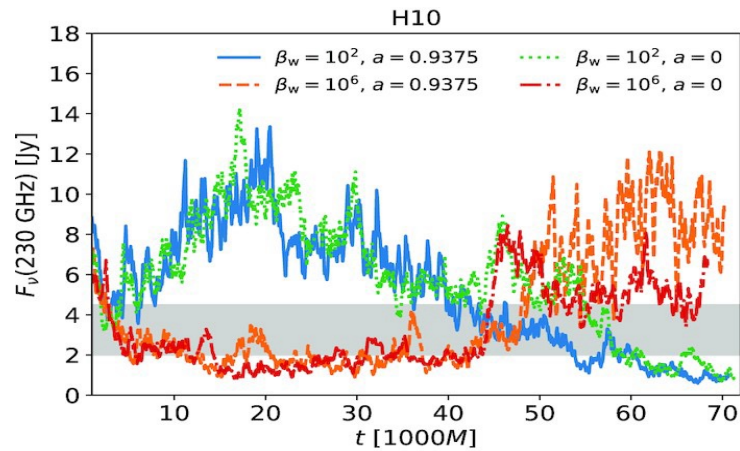
Electrons relativistic at \sim
100s of r_g

Low Electron Temperature
Near Horizon

Dimensionless Electron Temperature



230 GHz fluxes

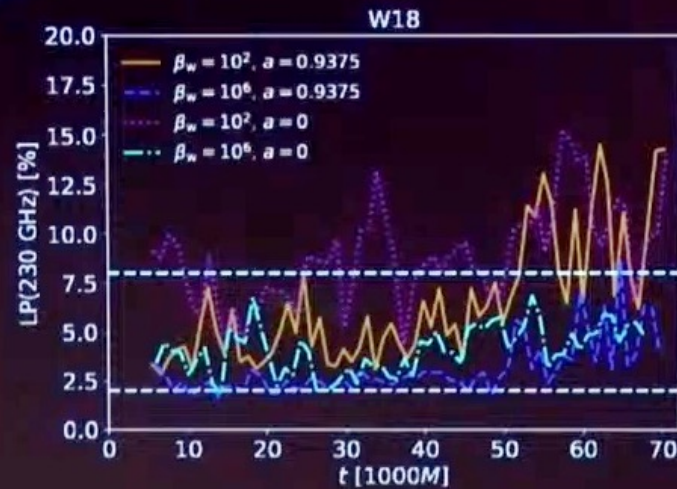
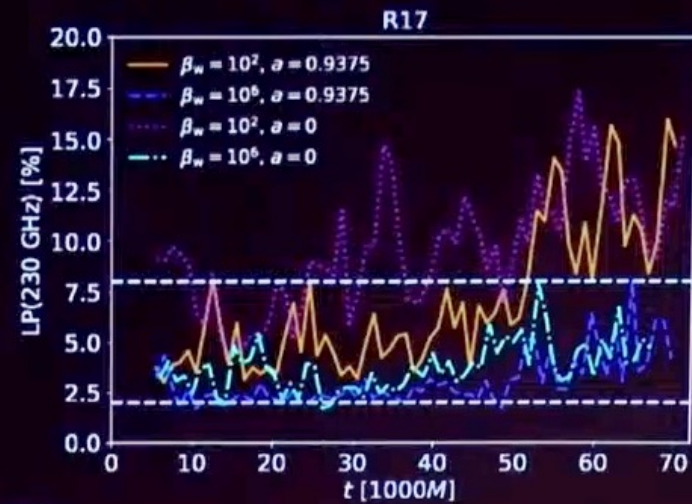
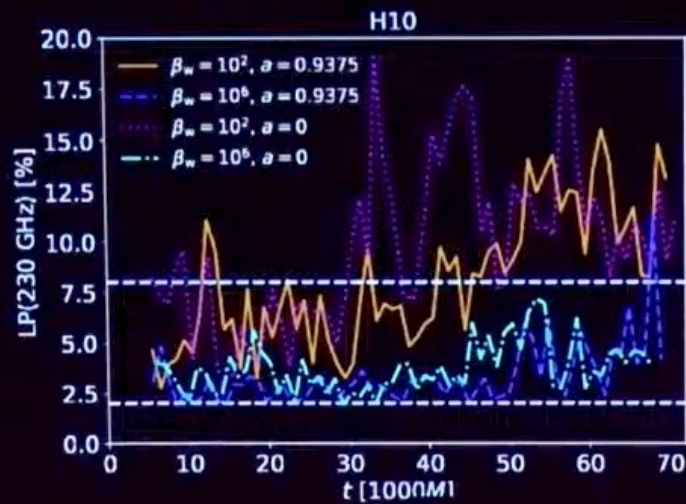


Shaded region: Observed value

Reconnection based heating
better!

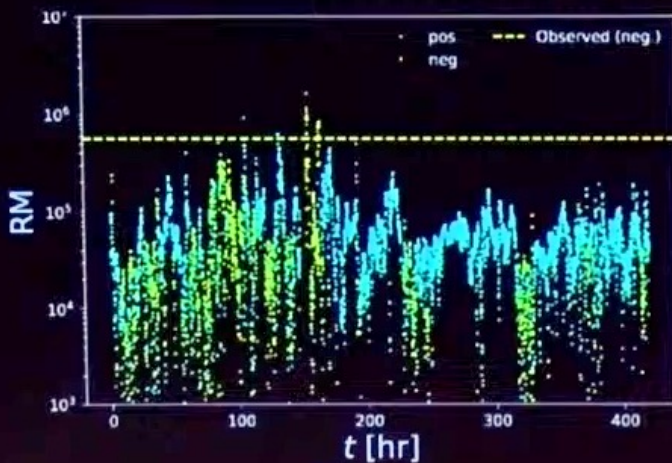
Weaker magnetic field case
more favourable!

Unresolved Linear Polarization Fractions



Rotation Measure

Small-Scale Flux Simulation



Observational Definition:

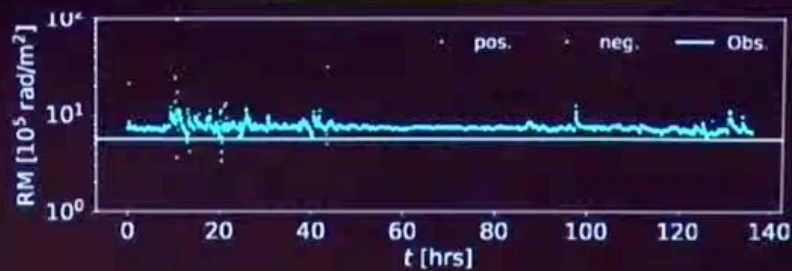
$$EVPA = EVPA_0 + RM\lambda^2$$

Theoretical Result for Relativistic Point Source

$$RM \propto \int f(\Theta_e) n_e \mathbf{B} \cdot d\mathbf{s}$$

$$f(\Theta_e) \sim 1/\Theta_e^2$$

Large-Scale Flux Simulation



Conclusions

- ❖ Sagittarius A* presents unique opportunity to study AGN accretion
 1. Bondi radius resolved
 2. Stellar winds that source accretion well constrained
- ❖ Wind-fed simulations can reasonably match observations (230 GHz flux, RM, LP fraction, X-ray luminosity, density profile)
- ❖ Strongly magnetized flows are predicted (though not necessarily MAD)
- ❖ Tilted disks/jets align during peak magnetic flux [i.e., when MAD or close to MAD]
- ❖ Observed RM requires consistent large-scale B-field

The trouble with people is not that they know so little, but that what they know is largely not true.

~ Mark Twain

Thank you.